

Numerical analysis of pressure drop in steady stenotic flows by using Lorentz's reciprocal theorem

ChangJin Ji^{1,*}, Kazuyasu Sugiyama², Shigeho Noda², Ying He¹ and Ryutaro Himeno²

¹ Department of Modern Mechanics, University of Science and Technology of China, China

² Advanced Center for Computing and Communication, RIKEN, Japan

* E-mail: jcj@mail.ustc.edu.cn

Key Words: Numerical simulation, Stenosis artery, VOF method, Lorentz's Reciprocal theorem, Lubrication theory.

The study of stenotic flow is very helpful in understanding the occurrence and development of many cardiovascular and cerebrovascular diseases, and thus is helpful in the diagnosis and treatment of these diseases. Our focus is on the issue how to systemically model the relation between the pressure drop and the flow rate in terms of the stenosis severity and the convective momentum transport.

Numerical simulations of simplified stenotic flows [Fig. 1] are performed by using a three-dimensional voxel-based simulator^[1], and the results based on Lorentz's reciprocal theorem^[2] and lubrication theory^[3] are analyzed. The effect of stenosis severities with a range of [10%, 80%] and Reynolds numbers with a range of [1-1000], which includes most of the arteries and arterioles in human body, on the pressure drop is considered. We derive an identity accounting for pressure drop mechanism, and demonstrate its validity for a computation of a steady flow.

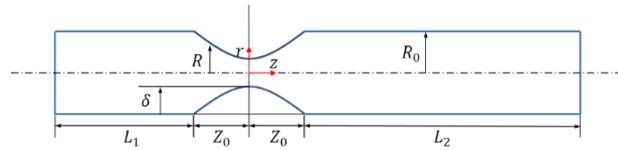


Fig.1 Simplified geometric model for stenosis vessel

Pressure drop in stenotic flows

If we call the solutions \mathbf{u} , p as the NS (Navier-Stokes) solutions and \mathbf{u}^* , p^* as the Stokes solutions, the Lorentz's reciprocal theorem^[4] can be represented as

$$(\nabla \mathbf{u}) : \boldsymbol{\sigma}^* = (\nabla \mathbf{u}^*) : \boldsymbol{\sigma}. \quad (1)$$

By jointly deducing the Eq.(1) and the continuity and momentum equations of two flows, we arrive at an identical relation

$$\underbrace{\overbrace{\Delta P^* Q}^{(W_1)} + \underbrace{\frac{2}{Re} \left(\iint_{\partial \Omega_{in} \& \partial \Omega_{out}} n_i u_j D_{ij}^* dS \right)}^{(W_2)}}_{(W_L)} = \underbrace{\overbrace{\Delta P Q^*}^{(W_3)} + \underbrace{\int (u_i u_j D_{ij}^*) dV}^{(W_4)} + \underbrace{\frac{2}{Re} \left(\iint_{\partial \Omega_{in} \& \partial \Omega_{out}} n_i u_j^* D_{ij} dS \right)}^{(W_5)}}_{(W_R)} - \underbrace{\left(\iint_{\partial \Omega_{in} \& \partial \Omega_{out}} n_i (u_i u_j u_j^*) dS \right)}_{(W_6)}. \quad (2)$$

Note that in the Stokes flow, the pressure drop ΔP^* is linear with respect to the flow rate Q^* . When ΔP^* is controlled to be $Q^* = Q$, the difference $\Delta P - \Delta P^*$ can be determined from the

spatial distribution of the velocity \mathbf{u}, \mathbf{u}^* as indicated in Eq. (2).

The computational results showed that the values of W_2 and W_5 , which are caused by the nonsymmetrical distribution of velocity at the two ends, are always much less than other terms, which is because that D_{ij}^* are equal at both sides due to the symmetry of Stokes flow and are almost 0 in the axis direction. While the difference of velocities is much less than Re , thus W_2 can be ignored. The W_6 appears to be small at low Reynolds number, but increases significantly with the increase of Reynolds number when Re is larger than 200. It is seen that W_6 is almost independent of the stenosis severity. Hence, when we set $Q=Q^*$, Eq. (2) can be simplified as

$$\Delta P = \Delta P^* + \underbrace{\frac{1}{Q} \int (-u_i u_j D_{ij}^*) dV}_{\Delta P_V} + \underbrace{\frac{1}{Q} \left(\iint n_i (u_i u_j u_j^*) dS \right)_{\partial\Omega_{in} \& \partial\Omega_{out}}}_{\Delta P_S}, \quad (3)$$

which means the pressure drop of NS flow through a stenosis tube can be decomposed into three parts. The first part is the pressure drop of Stokes flow through a stenosis tube and is mainly caused by viscous effect. The second part named as ΔP_V is the volume integral of $-u_i u_j D_{ij}^*$ within the whole fluid field and the third part named as ΔP_S is the integral of $u_i u_j u_j^*$ at two ends of tube. The last two terms are both caused by the uneven distributions of velocity on the both sides of stenosis region due to the presence of convective term in the Navier-Stokes Equations. The Eq. (3) offers a new method to analyze the pressure drop in stenotic flows, and the effects of these terms will be analyzed in details.

Stokes flow in stenosis tubes

The pressure drop of Stokes flow through a moderate stenosis can be approximately represented as

$$\Delta P^* = \frac{1}{Re} \frac{8Q}{\pi} \int_0^L R(z)^{-4} dz. \quad (4)$$

The variation of pressure drop obtained from the lubrication theory and Simulation results show a good agreement at mild and moderate stenosis, however there appears relatively obvious difference between them at higher stenosis severity.

Therefore, by adding a factor expressed as a function of stenosis severity St , the lubrication equation (12) can be modified as

$$\Delta P^* = \frac{1}{1-0.1St^2} \frac{1}{Re} \frac{8Q}{\pi} \int_0^L R(z)^{-4} dz. \quad (5)$$

The modification reduces the error of pressure drop produced by lubrication theory from 6% to below 0.1%.

REFERENCES

- [1] S. Noda, K. Fukasaku, and R. Himeno. Blood flow simulator using medical images without mesh generation, IFMBE Proc. of World Congress on Medical Physics and Biomedical Engineering, pp. 36-40, 2006.
- [2] J. Happel and H. Brenner. Low Reynolds Number Hydrodynamics. Martinus Nijhoff Publishers, The Hague, Chap. 3, 1983.
- [3] G. Batchelor. An introduction to fluid dynamics. Cambridge university press. Chap. 4, 2000.
- [4] M. Sbragaglia and K. Sugiyama. Boundary induced nonlinearities at small Reynolds numbers, Physica D, Vol. 228, pp.140-147, 2007.